

Fields with a Gaussian Measure

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In this paper, I calculate $\int \prod_{i=1}^n \delta(\Phi(t_i) - x_i) d\mu_c$. This corresponds to specifying the value of a field Φ defined by a gaussian measure $d\mu_c$ to take the values x_i at a set of n distinct times t_i .

In order to calculate $\int \prod_{i=1}^n \delta(\Phi(t_i) - x_i) d\mu_c$, I begin by calculating:

$$\int d\mu_c e^{i \int dv \Phi(v) g(v)} = \int d\mu_c \left\{ \sum_{n=0}^{\infty} \frac{i^n}{n!} \left(\int dv \Phi(v) g(v) \right)^n \right\}. \quad (1)$$

A generic term in the expansion of the exponential can be rewritten in terms multiple integration variables v_i :

$$\int d\mu_c \left(\int dv \Phi(v) g(v) \right)^n = \int d\mu_c \prod_{i=1}^n dv_i g(v_i) \Phi(v_i). \quad (2)$$

The integral over $d\mu_c$ will be non-zero only when n is even, in which case it takes the form:

$$\int d\mu_c \prod_{i=1}^n \Phi(v_i) = \sum_{\text{pairings}} C(v_{i_1} - v_{i_2}) \dots C(v_{i_{n-1}} - v_{i_n}) \quad (3)$$

where $C(v_i - v_j) = \frac{1}{2m} e^{-m|v_i - v_j|}$ and there are $(n-1)!!$ possible pairings. Since each term in the sum involves a product of $\frac{n}{2}$ two-point correlation functions and the $g(v_i)$ are symmetric under exchange of v_i :

$$\begin{aligned} & \int d\mu_c \prod_{i=1}^n dv_i g(v_i) \Phi(v_i) \\ &= (n-1)!! \left(\int dv_1 dv_2 g(v_1) g(v_2) C(v_1 - v_2) \right)^{\frac{n}{2}}. \end{aligned} \quad (4)$$

The original expression then simplifies to:

$$\begin{aligned} & \int d\mu_c e^{i \int dv \Phi(v) g(v)} \\ &= \sum_{n=0}^{\infty} \frac{i^{2n}}{2n!} (2n-1)!! \left(\int dv_1 dv_2 g(v_1) g(v_2) C(v_1 - v_2) \right)^n \\ &= e^{-\frac{1}{2} \int dv_1 dv_2 g(v_1) g(v_2) C(v_1 - v_2)}. \end{aligned} \quad (5)$$

Now, I can write $\int \prod_{i=1}^n \delta(\Phi(t_i) - x_i) d\mu_c$ in terms of

the Fourier transform of the delta function:

$$\begin{aligned} & \int \prod_{i=1}^n \delta(\Phi(t_i) - x_i) d\mu_c \\ &= \frac{1}{(2\pi)^n} \int d\mu_c \prod_{j=1}^n dk_j e^{ik_j(\Phi(t_j) - x_j)} \\ &= \frac{1}{(2\pi)^n} \int \prod_{j=1}^n dk_j e^{-i \sum_{j=1}^n k_j x_j} \int d\mu_c e^{i \int dv \Phi(v) [\sum_{j=1}^n k_j \delta(v - t_j)]}. \end{aligned} \quad (6)$$

Identifying $g(v) = \sum_{j=1}^n k_j \delta(v - t_j)$ and using Equation ?? gives:

$$\begin{aligned} & \int \prod_{i=1}^n \delta(\Phi(t_i) - x_i) d\mu_c \\ &= \frac{1}{(2\pi)^n} \int \prod_{j=1}^n dk_j e^{-i \sum_{j=1}^n k_j x_j} \\ & \quad \times e^{-\frac{1}{2} \int dv_1 dv_2 \sum_{j,\ell=1}^n k_j k_\ell \delta(v_1 - t_j) \delta(v_2 - t_\ell) C(v_1 - v_2)} \\ &= \frac{1}{(2\pi)^n} \int \prod_{j=1}^n dk_j e^{-i \sum_{j=1}^n k_j x_j} e^{-\frac{1}{2} \sum_{j,\ell=1}^n k_j k_\ell C(t_j - t_\ell)} \\ &= \frac{1}{(2\pi)^n} \int dk e^{-ik^\top x} e^{-\frac{1}{2} k^\top C k} \end{aligned} \quad (7)$$

where the last line is written in matrix form. Since $C_{ij} = C(t_i - t_j) = \frac{1}{2m} e^{-m|t_i - t_j|}$ is a real symmetric matrix, it can be diagonalized with an orthogonal matrix $P^{-1} = P^\top$ so that $PCP^\top = \Lambda$ and:

$$\begin{aligned} & \int \prod_{i=1}^n \delta(\Phi(t_i) - x_i) d\mu_c \\ &= \frac{1}{(2\pi)^n} \int dk e^{-ik^\top P^\top P x} e^{-\frac{1}{2} k^\top P^\top \Lambda P k} \\ &= \frac{1}{(2\pi)^n} \int \prod_{j=1}^n dk_j e^{-ik_j (Px)_j} e^{-\frac{1}{2} k_j^2 \Lambda_j} \\ &= \prod_{j=1}^n \frac{1}{\sqrt{2\pi \Lambda_j}} e^{-\frac{(Px)_j^2}{2\Lambda_j}} \end{aligned} \quad (8)$$

$$= \frac{1}{\sqrt{(2\pi)^n \det C}} e^{-\frac{1}{2} x^\top C^{-1} x}$$

When $n = 1$, $C(0) = \frac{1}{2m}$ and $\int \delta(\Phi(t) - x) d\mu_c = \sqrt{\frac{m}{\pi}} e^{-mx^2} = |\Omega_0(x)|^2$ the norm-squared of the ground state wave function for the quantum simple harmonic oscillator.